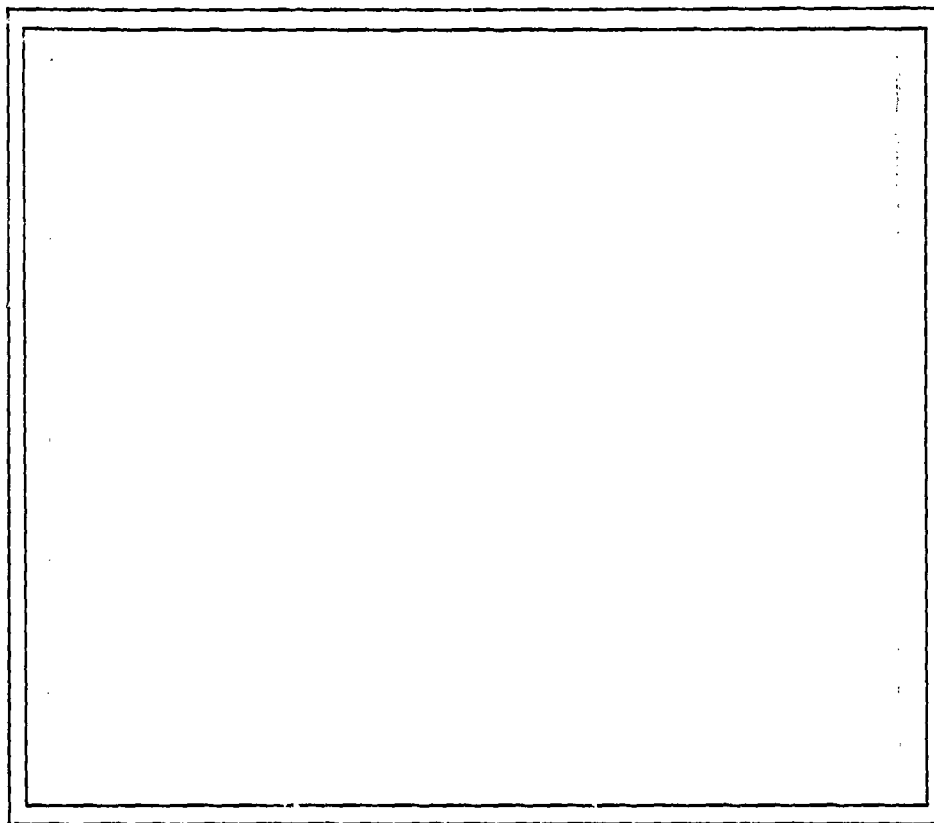


LEVEL II

①

ADA059118



DDC FILE COPY

PRINCETON UNIVERSITY  
Department of Civil Engineering

102 - 17.12.72



DDC  
RECEIVED  
SEP 26 1973  
D

DISTRIBUTION STATEMENT A

Approved for public release  
Distribution Unlimited

STRUCTURES AND MECHANICS

33

AD A059118

DDC FILE COPY

LEVEL II

①

Technical Rep. No. 53  
Civil Engng. Res. Rep. No. 78-SM-8

TR-53

14

⑥

LINE CRACK SUBJECT TO ANTIPLANE SHEAR

by

10

A. Cemal Eringen  
Princeton University

Research Sponsored by the  
Office of Naval Research

under

15

Contract, N00014-76-C-0240  
Modification No. P00002

⑨ Technical report

11 Jul 78

12

July 1978

Approved for public release; Distribution Unlimited

ADDITION for	
DTIC	White Section <input checked="" type="checkbox"/>
DDC	Buff Section <input type="checkbox"/>
ANNOUNCED	<input type="checkbox"/>
JUSTIFICATION	
DISTRIBUTION/AVAILABILITY CODES	

A

401 273

DDC  
RECEIVED  
SEP 26 1978  
REGULATED  
D

33

JOB

LINE CRACK SUBJECT TO  
ANTIPLANE SHEAR<sup>1</sup>

A. Cemal Eringen  
Princeton University

ABSTRACT

Field equations of nonlocal elasticity are solved to determine the state of stress in a plate with a line crack subject to a constant antiplane shear. Contrary to the classical elasticity solution, it is found that no stress singularity is present at the crack tip. By equating the maximum shear stress that occurs at the crack tip to the shear stress that is necessary to break the atomic bonds, the critical value of the applied shear is obtained for the initiation of fracture. If the concept of the surface tension is used, one is able to calculate the cohesive stress for brittle materials.

1. INTRODUCTION

In several previous papers [1] - [4] we discussed the state of stress near the tip of a sharp line crack in an elastic plate subject to uniform tension and in-plane shear. The field equations employed in the solution of these problems are those of the theory of nonlocal elasticity. The solutions obtained did not contain any stress singularity, thus resolving a fundamental problem that persisted over half a century. This enabled us to employ the maximum-stress hypothesis to deal with fracture problems in a natural way. Moreover it has been possible to predict the atomic cohesive stresses by introducing the experimental values of the surface energy.

---

<sup>1</sup>The present work was supported by the Office of Naval Research

The present paper deals with the problem of a line crack in an elastic plate where the crack surface is subject to a uniform anti-plane shear load. This problem, classically, is known as the Mode III displacement. We employ the field equations of nonlocal elasticity theory to formulate and solve this problem. The solution, as expected, does not contain the stress singularity at the crack tip and therefore a fracture criterion based on the maximum shear stress hypothesis can be used to obtain the critical value of the applied shear for which the line crack begins to become unstable. If the concept of the surface energy is introduced, it is possible to calculate the cohesive stress holding the atomic bonds together. Estimates of cohesive stress are also given for perfect crystalline solids.

In section 2 we present a resumé of basic equations of linear non-local elastic solids. In section 3 the boundary-value problem is formulated, and the general solution is obtained. In section 4 we give the solution of the dual integral equations, completing the solution. Calculations for the shear stress are carried out in a computer, and the results are discussed in section 5.

## 2. BASIC EQUATIONS OF NONLOCAL ELASTICITY

Basic equations of linear, homogeneous, isotropic, nonlocal elastic solids, with vanishing body and inertia forces, are (cf. [5]):

$$(2.1) \quad t_{kl,k} = 0$$

$$(2.2) \quad t_{kl} = \int_V [\lambda'(|\underline{x}' - \underline{x}|) e_{rr}(\underline{x}') \delta_{kl} + 2\mu'(|\underline{x}' - \underline{x}|) e_{kl}(\underline{x}')] dv(\underline{x}')$$

$$(2.3) \quad e_{kl} = \frac{1}{2}(u_{k,l} + u_{l,k})$$

where the only difference from classical elasticity is in the stress constitutive equations (2.2) in which the stress  $t_{k\ell}(\underline{x})$  at a point  $\underline{x}$  depends on the strains  $e_{k\ell}(\underline{x}')$ , at all points of the body. For homogeneous and isotropic solids there exist only two material moduli,  $\lambda'(|\underline{x}'-\underline{x}|)$  and  $\mu'(|\underline{x}'-\underline{x}|)$  which are functions of the distance  $|\underline{x}'-\underline{x}|$ . The integral in (2.2) is over the volume  $V$  of the body enclosed within a surface  $\partial V$ .

Throughout this paper we employ cartesian coordinates  $x_k$  with the usual convention that a free index takes the values (1, 2, 3), and repeated indices are summed over the range (1, 2, 3). Indices following a comma represent partial differentiation, e.g.

$$u_{k,\ell} = \partial u_k / \partial x_\ell$$

In our previous work [6, 7] we obtained the forms of  $\lambda'(|\underline{x}'-\underline{x}|)$  and  $\mu'(|\underline{x}'-\underline{x}|)$  for which the dispersion curves of plane elastic waves coincide with those known in lattice dynamics. Among several possible curves the following has been found very useful

$$\begin{aligned} (\lambda', \mu') &= (\lambda, \mu) \alpha(|\underline{x}'-\underline{x}|) \quad , \\ (2.4) \quad \alpha(|\underline{x}'-\underline{x}|) &= \alpha_0 \exp[-(\beta/a)^2 (\underline{x}'-\underline{x}) \cdot (\underline{x}'-\underline{x})] \quad , \end{aligned}$$

where  $\beta$  is a constant,  $a$  is the lattice parameter, and  $\alpha_0$  is determined by the normalization

$$(2.5) \quad \int_V \alpha(|\underline{x}'-\underline{x}|) dv(\underline{x}') = 1 \quad .$$

In the present work we employ the nonlocal elastic moduli given by (2.4)<sub>2</sub>. Carrying (2.4)<sub>2</sub> into (2.5) we obtain

$$(2.6) \quad \alpha_0 = \frac{1}{\pi} (\beta/a)^2 .$$

Substituting (2.4)<sub>1</sub> into (2.2) we write

$$(2.7) \quad t_{kl} = \int_V \alpha(|\underline{x}' - \underline{x}|) \sigma_{kl}(\underline{x}') dv(\underline{x}') ,$$

where

$$(2.8) \quad \begin{aligned} \sigma_{kl}(\underline{x}') &= \lambda e_{rr}(\underline{x}') \delta_{kl} + 2\mu e_{kl}(\underline{x}') \\ &= \lambda u_{r,r}(\underline{x}') \delta_{kl} + \mu [u_{k,l}(\underline{x}') + u_{l,k}(\underline{x}')] \end{aligned}$$

is the classical Hooke's law. Substituting (2.8) into (2.1) and using Green-Gauss theorem we obtain:

$$(2.9) \quad \int_V \alpha(|\underline{x}' - \underline{x}|) \sigma_{kl,k}(\underline{x}') dv(\underline{x}') - \oint_{\partial V} \alpha(|\underline{x}' - \underline{x}|) \sigma_{kl}(\underline{x}') da_k(\underline{x}') = 0 .$$

The contribution to the surface integral from the parts of the surface at infinity would be dropped since the displacement field vanishes at infinity.

### 3. CRACK UNDER ANTIPLANE SHEAR

We consider an elastic plate in the  $(x_1=x, x_2=y)$  - plane weakened by a line crack of length  $2l$  along the  $x$ -axis. The plate is subjected to a constant anti-plane shear stress  $t_{yz}=\tau_0$  along the surfaces of the crack, Fig. 1. For this problem we have

$$(3.1) \quad u_1=u_2=0, \quad u_3=w(x,y),$$

$$(3.2) \quad \sigma_{xz}=\mu \frac{\partial w}{\partial x}, \quad \sigma_{yz}=\mu \frac{\partial w}{\partial y}, \quad \text{all other } \sigma_{kl}=0,$$

so that the only surviving member of the field equations (2.9) is

$$(3.3) \quad \mu \oint \alpha(|x'-x|, |y'-y|) \nabla'^2 w(x', y') d\lambda' dy' - \int_{-l}^l \alpha(|x'-x|, |y|) [\sigma_{yz}(x', 0)] dx' = 0,$$

where the integral with a slash is over the two-dimensional infinite space excluding the line of the crack  $(|x| < l, y=0)$ . A boldface bracket indicates a jump at the crack line.

When an undeformed and unstressed body is sliced to create a free surface, it will in general be deformed and stressed on account of the long-range interatomic forces. Thus if we are to consider that the plate with a crack is undeformed and unstressed in its natural state then we must apply the boundary conditions on the unopened crack surface.

Under the applied anti-plane shear load on the unopened surfaces of the crack, the displacement field possesses the following symmetry regulations

$$(3.4) \quad w(x, -y) = -w(x, y).$$

Using this in (3.2) we find that

$$(3.5) \quad [\bar{\sigma}_{yz}(x,0)] = 0 \quad .$$

Hence the line integral in (3.3) vanishes. By taking the Fourier transform of (3.3) with respect to  $x'$ , we can show that the general solution of (3.3) is identical to that of

$$(3.6) \quad \frac{d^2 \bar{w}(\xi, y)}{dy^2} - \xi^2 \bar{w}(\xi, y) = 0 \quad ,$$

almost everywhere. Here a superposed bar indicates the Fourier transform e.g.

$$\bar{f}(\xi, y) = (2\pi)^{-\frac{1}{2}} \int_{-\infty}^{\infty} f(x, y) \exp(i\xi x) dx \quad .$$

The boundary conditions are

$$(3.7) \quad \begin{aligned} w(x,0) &= 0 & \text{for } |x| > l & , \\ \tau_{yz}(x,0) &= \tau_0 & \text{for } |x| < l & , \\ w(x,y) &= 0 & \text{as } (x^2 + y^2)^{1/2} \rightarrow \infty & . \end{aligned}$$

The general solution of (3.6) (for  $y \geq 0$ ) satisfying (3.7)<sub>3</sub> is

$$(3.8) \quad w(x,y) = (2/\pi)^{\frac{1}{2}} \int_0^{\infty} A(\xi) e^{-\xi y} \cos(\xi x) d\xi \quad ,$$

where  $A(\xi)$  is to be determined from the remaining two boundary conditions.

For the non-zero components of the stress tensor we have



$$\begin{aligned}
 t_{xz} &= -(2/\pi)^{\frac{1}{2}} \mu \int_0^\infty A(\xi) \xi d\xi \int_0^\infty dy' \int_{-\infty}^\infty [ \alpha(|x'-x|, |y'-y|) \\
 &\quad - \alpha(|x'-x|, |y'+y|) ] e^{-\xi y'} \sin(\xi x') dx', \\
 (3.9) \quad t_{yz} &= -(2/\pi)^{\frac{1}{2}} \mu \int_0^\infty A(\xi) \xi d\xi \int_0^\infty dy' \int_{-\infty}^\infty [ \alpha(|x'-x|, |y'-y|) \\
 &\quad + \alpha(|x'-x|, |y'+y|) ] e^{-\xi y'} \cos(\xi x') dx'.
 \end{aligned}$$

Using (2.4)<sub>2</sub> for  $\alpha(|x'-x|, |y'-y|)$ , we carry out integrations on  $x'$  and  $y'$ . To this end we note the following integrals, [8]:

$$\begin{aligned}
 I_1 &= \int_{-\infty}^\infty \exp(-px'^2) \left\{ \frac{\sin \xi(x'+x)}{\cos \xi(x'+x)} \right\} dx' = (\pi/p)^{\frac{1}{2}} \exp(-\xi^2/4p) \left\{ \frac{\sin(\xi x)}{\cos(\xi x)} \right\}, \\
 (3.10) \quad I_2 &= \int_0^\infty \exp(-py'^2 - \gamma y') dy' = \frac{1}{2}(\pi/p)^{\frac{1}{2}} \exp(\gamma^2/4p) [1 - \Phi(\gamma/2\sqrt{p})], \\
 \Phi(z) &\equiv 2\pi^{-\frac{1}{2}} \int_0^z \exp(-t^2) dt.
 \end{aligned}$$

Hence

$$\begin{aligned}
 t_{xz} &= -(2\pi)^{-\frac{1}{2}} \mu \int_0^\infty \xi A(\xi) [e^{-\xi y} \operatorname{erfc}\left(\frac{\xi-2py}{2\sqrt{p}}\right) - e^{\xi y} \operatorname{erfc}\left(\frac{\xi+2py}{2\sqrt{p}}\right)] \sin(\xi x) d\xi, \\
 (3.11) \quad t_{yz} &= -(2\pi)^{-\frac{1}{2}} \mu \int_0^\infty \xi A(\xi) [e^{-\xi y} \operatorname{erfc}\left(\frac{\xi-2py}{2\sqrt{p}}\right) + e^{\xi y} \operatorname{erfc}\left(\frac{\xi+2py}{2\sqrt{p}}\right)] \cos(\xi x) d\xi, \\
 p &\equiv (\beta/a)^2, \quad \operatorname{erfc}(z) = 1 - \Phi(z).
 \end{aligned}$$

The boundary conditions (3.7)<sub>1</sub> and (3.7)<sub>2</sub> now read

$$\begin{aligned}
 (3.12) \quad & \int_0^{\infty} \zeta^{\frac{1}{2}} C(\zeta) K(\epsilon \zeta) \cos(z\zeta) d\zeta = -(\pi/2)^{\frac{1}{2}} T_0, & 0 < z < 1, \\
 & \int_0^{\infty} \zeta^{-\frac{1}{2}} C(\zeta) \cos(z\zeta) d\zeta = 0, & z > 1,
 \end{aligned}$$

where we set

$$\begin{aligned}
 (3.13) \quad & z = x/l, \quad \zeta = \xi l, \quad \epsilon = a/2\beta l, \\
 & K(\epsilon \zeta) = \operatorname{erfc}(\epsilon \zeta), \\
 & A(\xi) = \zeta^{-\frac{1}{2}} C(\zeta), \quad T_0 = \tau_0 l^2/\mu.
 \end{aligned}$$

To determine the unknown function  $A(\xi)$ , we must solve the dual integral equations (3.12).

#### 4. THE SOLUTION OF THE DUAL INTEGRAL EQUATIONS

Recalling the expression

$$\cos(z\zeta) = (\pi z\zeta/2)^{\frac{1}{2}} J_{-\frac{1}{2}}(z\zeta),$$

where  $J_\nu(z)$  is the Bessel function of order  $\nu$ , we write the system (3.12) in the form

$$\begin{aligned}
 (4.1) \quad & \int_0^{\infty} \zeta C(\zeta) [1+k(\epsilon \zeta)] J_{-\frac{1}{2}}(z\zeta) d\zeta = -T_0 z^{-\frac{1}{2}}, & 0 < z < 1, \\
 & \int_0^{\infty} C(\zeta) J_{-\frac{1}{2}}(z\zeta) d\zeta = 0, & z > 1.
 \end{aligned}$$

The kernel function  $k(\epsilon\zeta)$  is given by

$$(4.2) \quad k(\epsilon\zeta) = K(\epsilon\zeta) - 1 = -\phi(\epsilon\zeta) .$$

The solution of the dual integral equations (4.1) is not known. However, it is possible to reduce the problem to the solution of a Fredholm equation (cf. [9])

$$(4.3) \quad h(x) + \int_0^1 h(u) L(x, u) du = -\frac{1}{2}(\pi x)^{\frac{1}{2}} T_0 ,$$

for the function  $h(x)$ , where

$$(4.4) \quad L(x, u) = (xu)^{\frac{1}{2}} \int_0^{\infty} t k(\epsilon t) J_0(xt) J_0(ut) dt .$$

When (4.3) is solved, then  $C(\zeta)$  is calculated by

$$(4.5) \quad C(\zeta) = (2\zeta)^{\frac{1}{2}} \int_0^1 x^{\frac{1}{2}} J_0(\zeta x) h(x) dx .$$

As discussed in a previous work [4], if we note that  $\epsilon$  is extremely small,  $k(\epsilon\zeta)$  may be neglected as compared to unity in (4.1), (see Fig. 2). In this case the zeroth order solution of (4.3) namely  $h_0(x) = -T_0(\pi x)^{\frac{1}{2}}/2$  suffices for the calculations when the crack size is larger than 100 atomic distances. In such a case we have

$$(4.6) \quad C_0(\zeta) = -(\pi/2)^{\frac{1}{2}} T_0 \zeta^{-\frac{1}{2}} J_1(\zeta) ,$$

and therefore

$$(4.7) \quad A_0(\xi) = -(\pi/2)^{1/2} T_0 J_1(\xi \ell) / \xi \ell$$

The shear stresses are then calculated by (3.11). Interesting among these is the shear stress  $t_{yz}$  along the crack line  $y = 0$ . For this we obtain

$$(4.8) \quad t_{yz}(z, 0) / \tau_0 = \int_0^\infty K(\epsilon \zeta) J_1(\zeta) \cos(z\zeta) d\zeta$$

As observed before, this integral converges for all  $z$  provided  $K(\epsilon \zeta)$  is not approximated by unity for  $\epsilon$  small. For  $\epsilon = 0$  at  $z = 1$  we have the classical stress singularity. However, so long as  $\epsilon \neq 0$ , (4.8) gives a finite stress all along  $y = 0$ . At  $0 < z < 1$ ,  $t_{yz} / \tau_0$  is very close to unity, and for  $z > 1$ ,  $t_{yz} / \tau_0$  possesses finite values diminishing from a maximum value at  $z=1$  to zero at  $z=\infty$ .

For  $\epsilon \gg 1/100$  the approximate solution given by (4.7) is not very good. However, further improvements can be achieved by the iterative solution of (4.3) with the use of  $C_0(\zeta)$ . Since  $\epsilon \gg 1/100$  represents a crack length of less than  $10^{-6}$  cm, and at such submicroscopic sizes other serious questions arise regarding the interatomic arrangements and force laws, we do not pursue solutions valid at such small crack sizes.

##### 5. NUMERICAL CALCULATIONS AND DISCUSSION.

Calculations of the shear stress  $t_{yz}$ , given by (4.8) along the crack line, were carried out on a computer. The results are plotted for  $\epsilon = 1/20, 1/50, 1/100, 1/200$ , in Figures 3 to 6. For a crack length of 20 atomic distances ( $\epsilon = 1/20$ ) the result is not very good in that the boundary

condition at  $|x| < l$ ,  $y = 0$  is satisfied only very roughly. However, for a crack size of 100 atomic distances (Fig. 5) the shear stress boundary condition is fulfilled in a strong approximate sense. The relative error in this case is less than  $1\frac{1}{2}\%$ . Hence we conclude that the classical  $A_0(\xi)$  given by (4.7) gives satisfactory results for crack lengths greater than 100 atomic distances.

The stress concentration occurs at the crack tip, and this is given by

$$(5.1) \quad t_{yz}(l,0)/t_0 = c_3/\sqrt{\epsilon} \quad , \quad \epsilon \equiv a/2\beta l \quad ,$$

where  $c_3$  converges to about

$$(5.2) \quad c_3 \approx 0.40 \quad .$$

The following observations are very significant:

- (i) The maximum shear stress occurs at the crack tip, and it is finite (eq. 5.1)
- (ii) The shear stress at the crack tip becomes infinite as the atomic distance  $a \rightarrow 0$ . This is the classical continuum limit of square root singularity.
- (iii) When  $t_{yz}(l,0) = t_c$  (= cohesive shear stress), the plate will fail.  
In this case

$$(5.3) \quad \tau_0^2 l = C_G$$

where

$$(5.4) \quad C_G = (a/2\beta c_3^2) t_c^2$$

Equation (5.3) is the expression of the Griffith fracture criterion for brittle fracture. We have arrived at this result via the maximum shear-stress hypothesis, rather than the surface energy concept used by Griffith and his followers. The significance of this result is that the fracture criteria are unified at both the macroscopic and the microscopic scales and that the natural concept of bond failure is employed.

(iv) The cohesive shear stress  $t_c$  may be estimated if one employs the Griffith's definition of the surface energy  $\gamma$  and writes

$$(5.5) \quad t_c^2 a = K_c \gamma ,$$

where

$$(5.6) \quad K_c = 8\mu c_3^2 \beta / \pi(1-\nu)$$

Since some measurements exist on  $\gamma$ , by employing these values we can calculate the cohesive shear stress. For steel we have

$$\gamma = 1975 \text{ CGS} , \quad \mu = 6.92 \times 10^{11} \text{ CGS}$$

$$\nu = 0.291 , \quad a = 2.48 \text{ \AA}^2$$

$$(5.7) \quad t_c / \mu = 0.2568 \beta^{1/2}$$

Let at a  $n$  atomic distance the nonlocal effects attenuate to 1% of its value at  $x=0$ . Using (2.4) we find that

$$(5.8) \quad \beta = 2.146/n ,$$

$$(5.9) \quad t_c/\mu = 0.3764 n^{-1/2} .$$

For  $n = 6$  this gives

$$(5.10) \quad t_c/\mu = 0.14$$

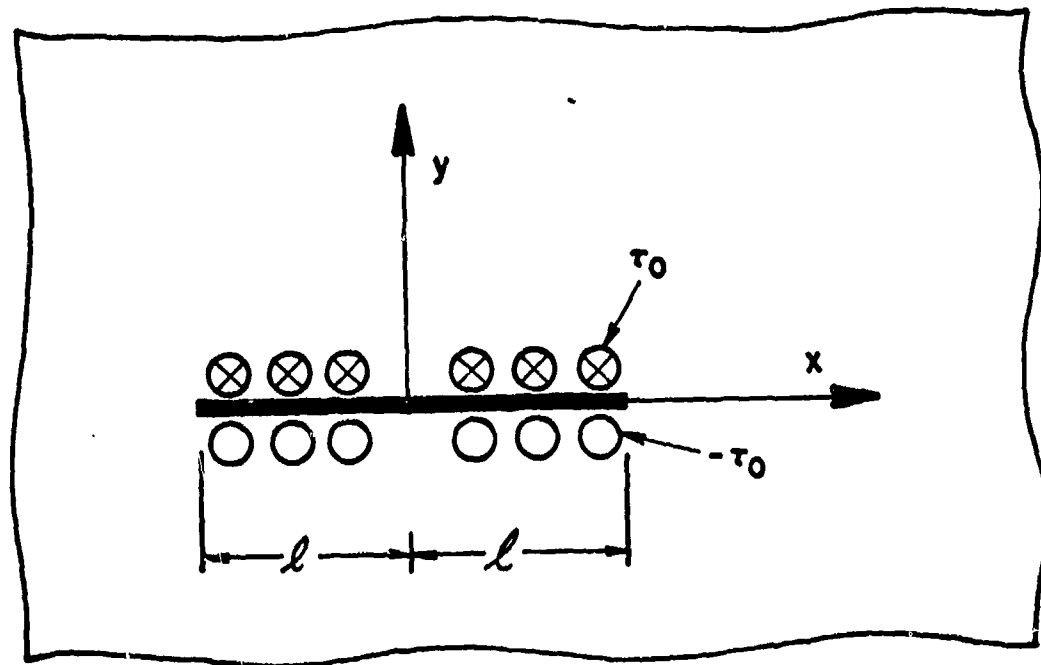
which is in the right range and well accepted by metallurgists. For example Kelly [10] gives  $t_c/\mu = 0.11$ . There is however a question of in-plane versus antiplane shear failure which need special examination. Compared to the results obtained in our previous work on the in-plane shear [4], the cohesive stress seems 30% higher. However, this is somewhat artificial since it is necessary to know the value of  $\beta$  or  $n$  in either case more precisely. This of course requires at least one experiment.

Acknowledgement: The author is indebted to Dr. Balta for carrying out computations.

# REFERENCES

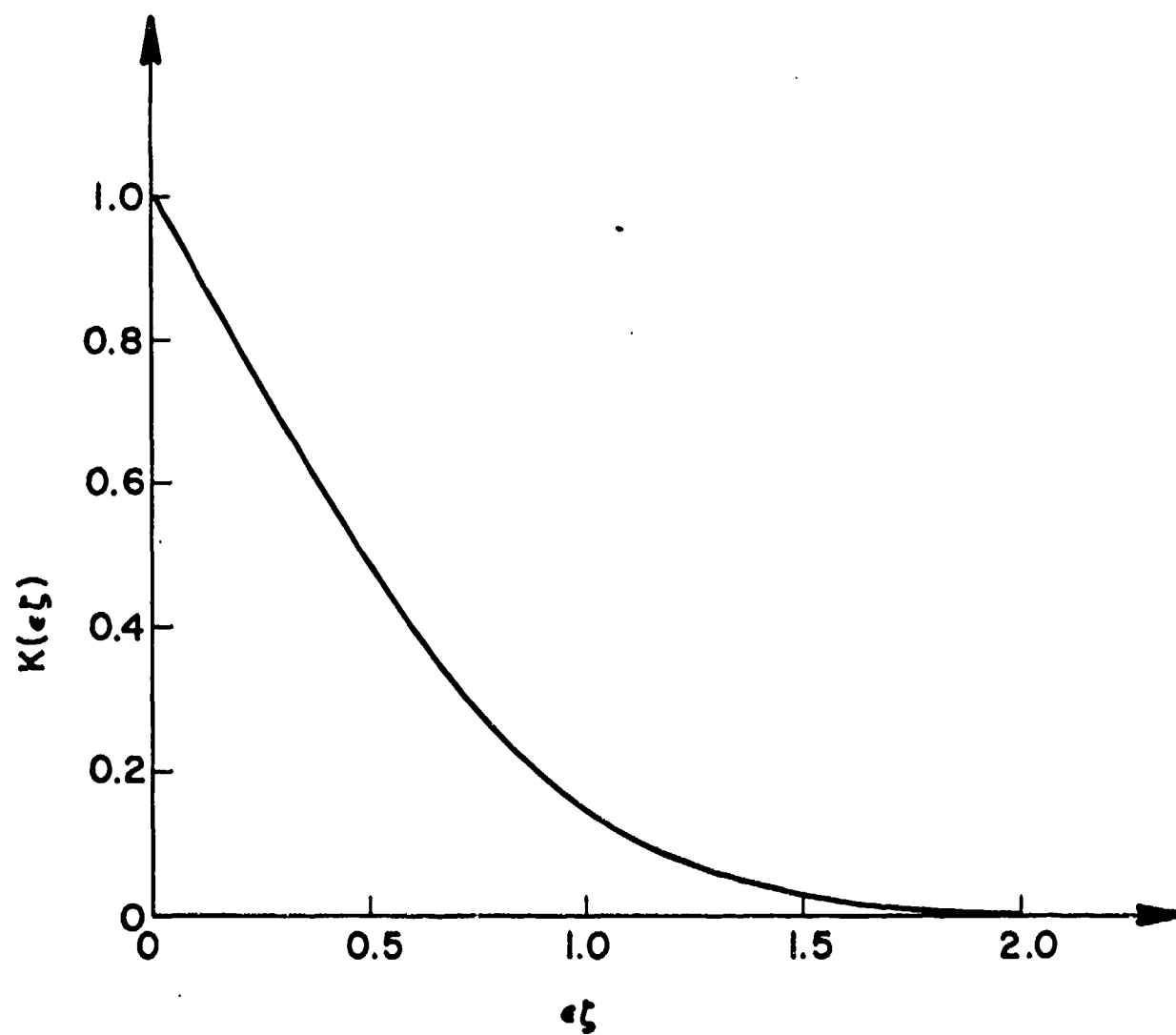
- [1] A. C. Eringen and B. S. Kim, "Stress Concentration at the Tip of Crack," Mech. Res. Comm. 1, 233-237, 1974.
- [2] A. C. Eringen, "State of Stress in the Neighborhood of a Sharp Crack Tip," Trans. of the twenty-second conference of Army mathematicians, 1-18, 1977.
- [3] A. C. Eringen, C. G. Speziale and B. S. Kim, "Crack Tip Problem in Nonlocal Elasticity", J. Mech. and Phys. of Solids, 25, 339-355, 1977.
- [4] A. C. Eringen, "Line Crack Subject to Shear", to appear in Int. J. Fracture.
- [5] A. C. Eringen, "Linear Theory of Nonlocal Elasticity and Dispersion of Plane Waves", Int. J. Engng. Sci. 10, 233-248, 1972.
- [6] A. C. Eringen, "Nonlocal Elasticity and Waves," Continuum Mechanics Aspects of Geodynamics and Rock Fracture Mechanics, (ed. P. Thoft-Christensen) Dordrecht, Holland, D. Reidel Publishing Co., pp. 81-105, 1974
- [7] A. C. Eringen, "Continuum Mechanics at the Atomic Scale", Cryst. Lattice Defects, 7, pp. 109-130, 1977.
- [8] I. S. Gradshteyn and I. W. Ryzhik, "Tables of Integrals, Series and Products", New York; Acad. Press, 1965, pp. 480, 307, 338.
- [9] I.N. Sneddon, "Mixed Boundary Value Problems in Potential Theory", North Holland Publishing Co., Amsterdam, 1966, pp. 106-108.
- [10] A. Kelley, "Strong Solids", Oxford, 1966, p. 19.





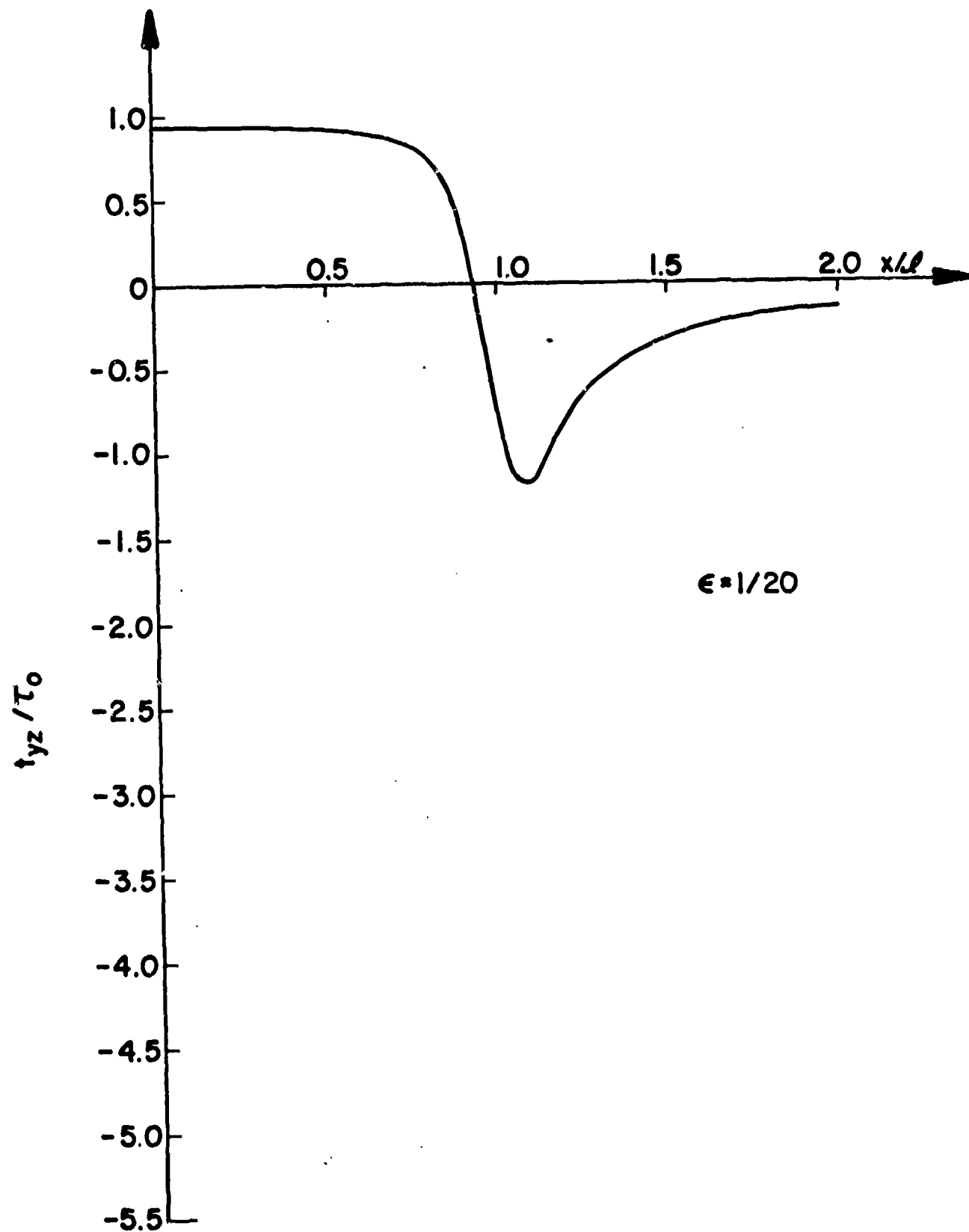
LINE CRACK SUBJECT TO ANTIPLANE SHEAR

FIGURE 1



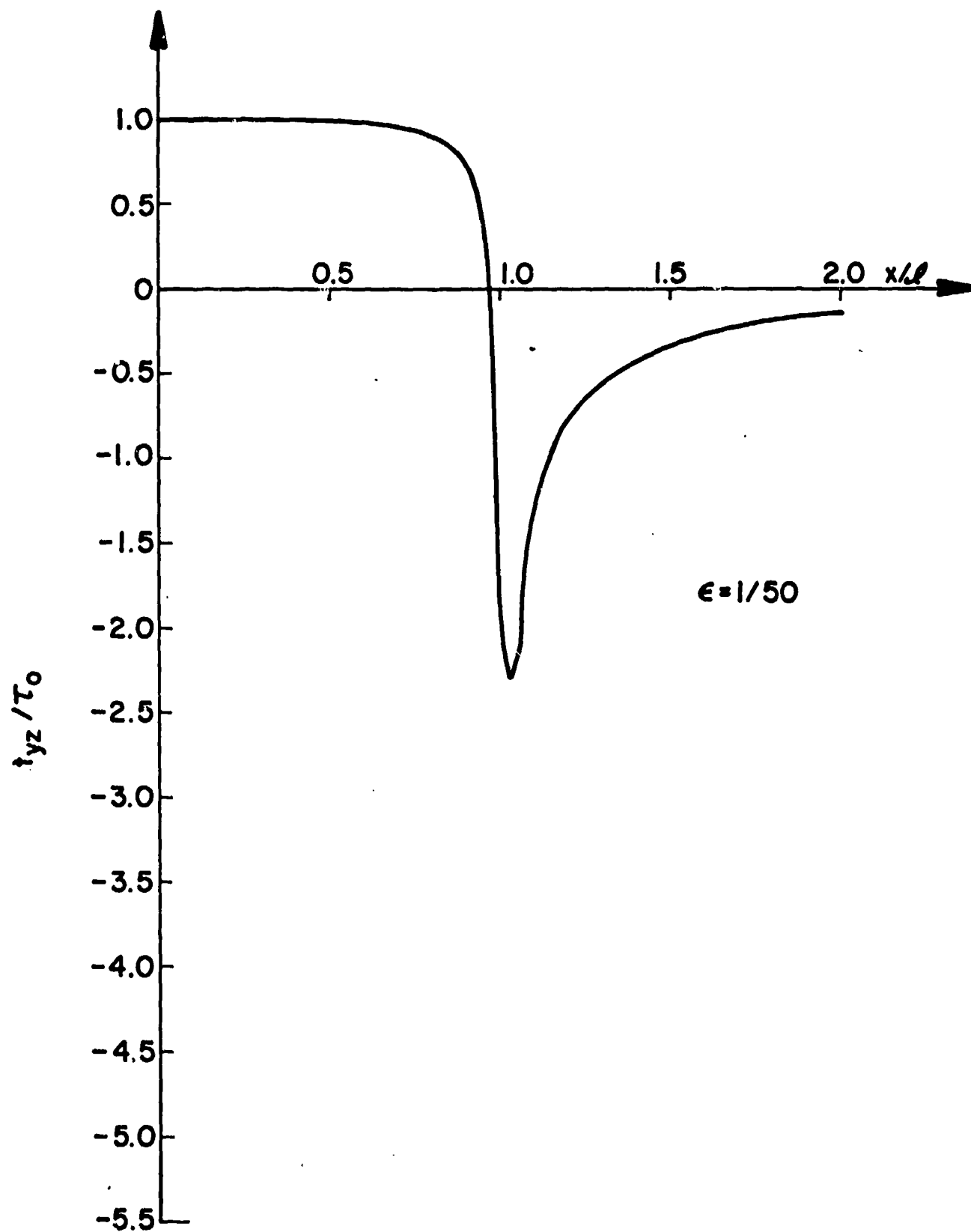
KERNEL FUNCTION  $K(\epsilon\zeta)$

FIGURE 2



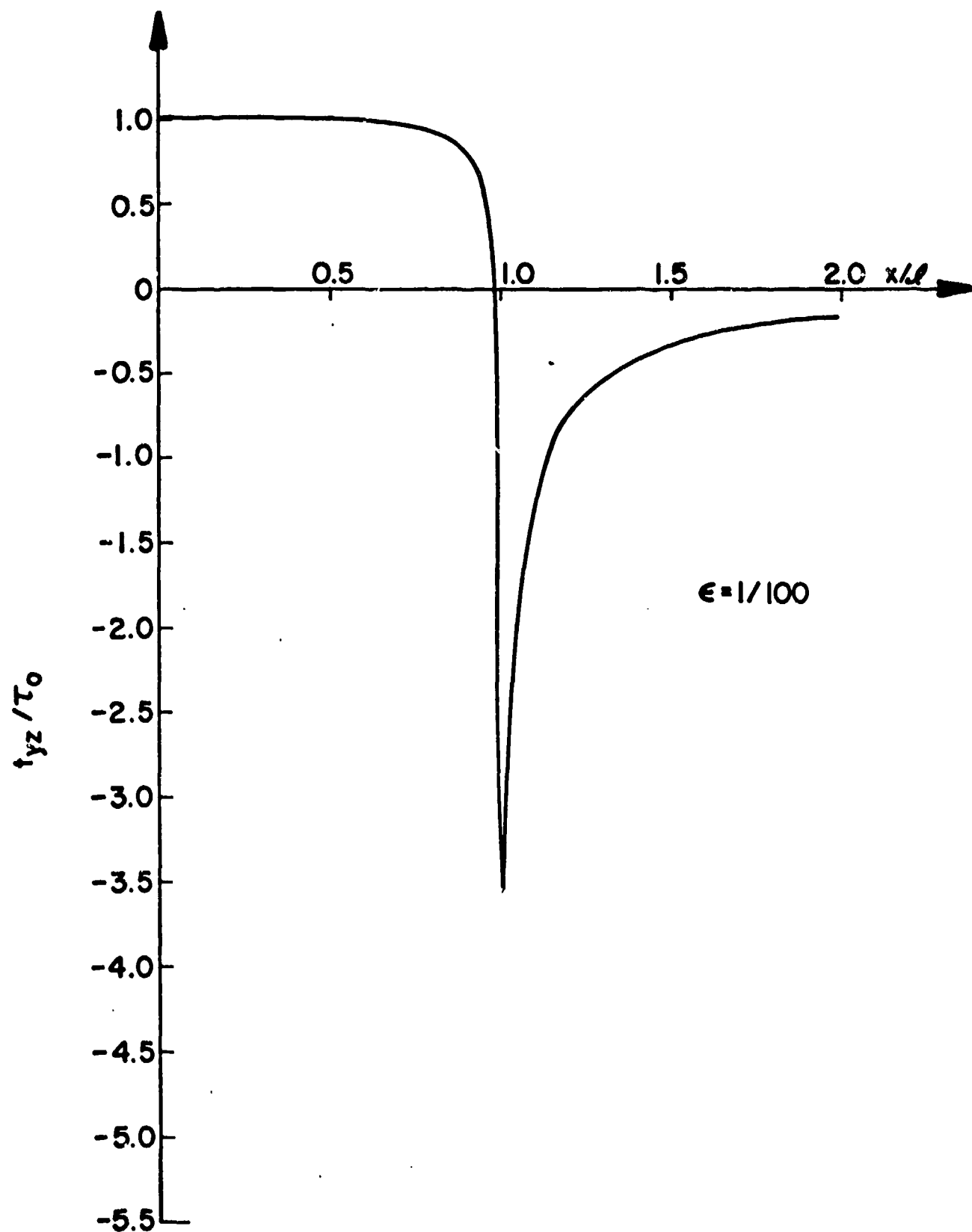
ANTIPLANE SHEAR STRESS

FIGURE 3



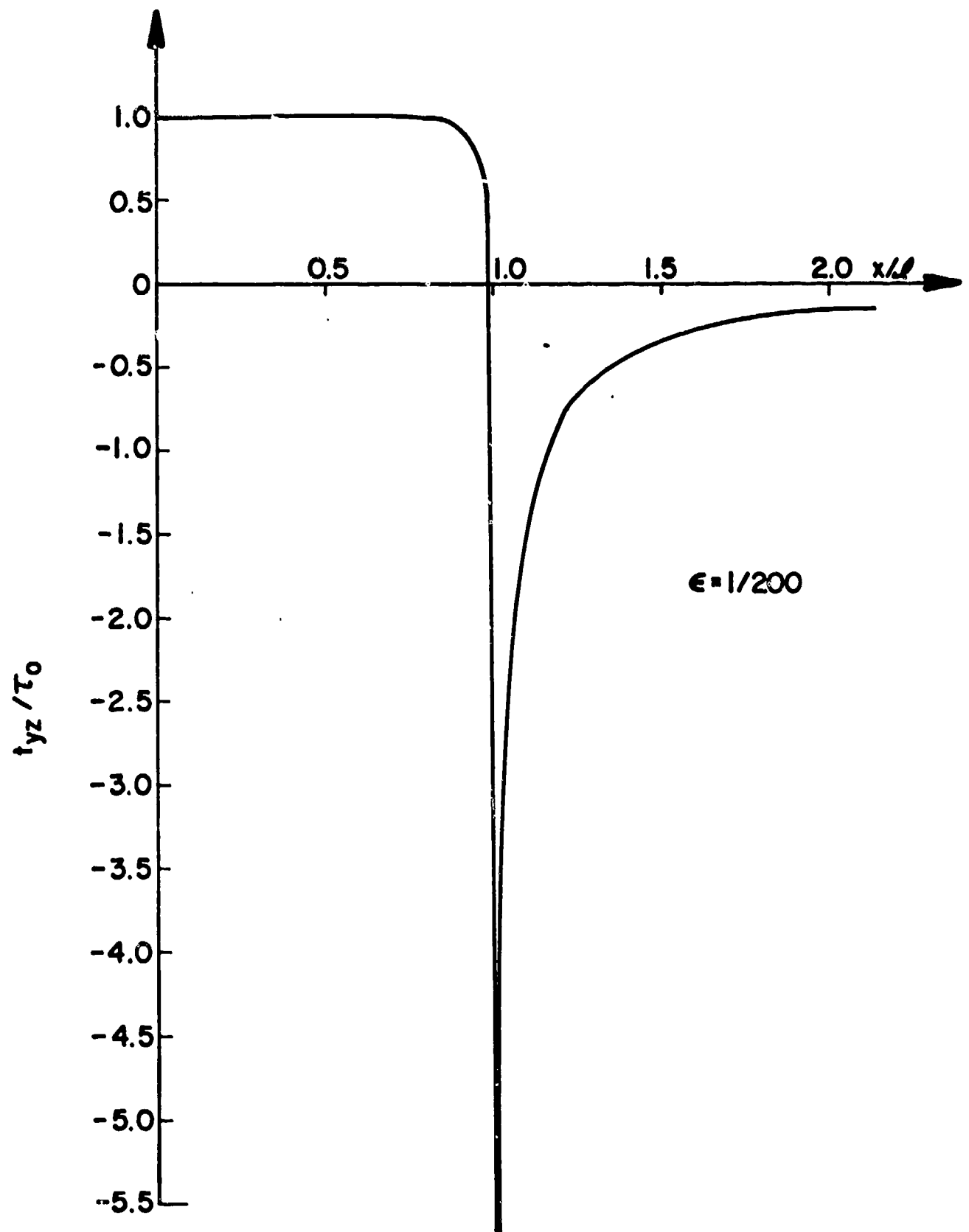
ANTIPLANE SHEAR STRESS

FIGURE 4



ANTIPLANE SHEAR STRESS

FIGURE 5



ANTIPLANE SHEAR STRESS

FIGURE 6

REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER Princeton Technical Rep. 53	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER
4. TITLE (and Subtitle) Line Crack Subject to Antiplane Shear		5. TYPE OF REPORT & PERIOD COVERED Technical Report
		6. PERFORMING ORG. REPORT NUMBER
7. AUTHOR(s) A. Cemal Eringen		8. CONTRACT OR GRANT NUMBER(s) N00014-76-C-0240
9. PERFORMING ORGANIZATION NAME AND ADDRESS Princeton University Princeton, NJ 08540		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS P00002
11. CONTROLLING OFFICE NAME AND ADDRESS Office of Naval Research (Code 471) Arlington, VA 22217		12. REPORT DATE July, 1978
		13. NUMBER OF PAGES 14 Pages
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office)		15. SECURITY CLASS. (of this report) Unclassified
		15a. DECLASSIFICATION/DOWNGRADING SCHEDULE
16. DISTRIBUTION STATEMENT (of this Report) Approved for public release; Distribution Unlimited		
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)		
18. SUPPLEMENTARY NOTES		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) fracture mechanics, line crack, antiplane shear, nonlocal theory, crack tip problems, cohesive stress		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) Field equations of nonlocal elasticity are solved to determine the state of stress in a plate with a line crack subject to a constant anti-plane shear. Contrary to the classical elasticity solution, it is found that no stress singularity is present at the crack tip. By equating the maximum shear stress that occurs at the crack tip to the shear stress that is necessary to break the atomic bonds, the critical value of the applied shear is obtained for the initiation of fracture. If the concept of the surface tension is used, one is able to calculate the cohesive stress for brittle materials.		

DD FORM 1473

1 JAN 73

EDITION OF 1 NOV 68 IS OBSOLETE  
S/N 0102-014-6601

SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

PART 1 - GOVERNMENT

Administrative & Liaison Activities

Chief of Naval Research  
Department of the Navy  
Arlington, Virginia 22217

Attn: Code 474 (2)

471

222

Director  
ONR Branch Office  
495 Summer Street  
Boston, Massachusetts 02210

Director  
ONR Branch Office  
219 S. Dearborn Street  
Chicago, Illinois 60604

U.S. Naval Research Laboratory  
Attn: Code 2627  
Washington, D.C. 20390

Director  
ONR - New York Area Office  
715 Broadway - 5th Floor  
New York, N.Y. 10003

Director  
ONR Branch Office  
1030 E. Green Street  
Pasadena, California 91101

Defense Documentation Center  
Cameron Station  
Alexandria, Virginia 22314 (12)

Army

Commanding Officer  
U.S. Army Research Office Durham  
Attn: Mr. J. J. Murray  
CRD-AA-IP  
Box CM, Duke Station  
Durham, North Carolina 27706 . 2 .

Commanding Officer

AMXMR-ATL

Attn: Mr. R. Shea  
U.S. Army Materials Res. Agency  
Watertown, Massachusetts 02172

Watervliet Arsenal

MAGGS Research Center

Watervliet, New York 12189

Attn: Director of Research

Technical Library

Redstone Scientific Info. Center  
Chief, Document Section  
U.S. Army Missile Command  
Redstone Arsenal, Alabama 35809

Army R&D Center  
Fort Belvoir, Virginia 22060

Navy

Commanding Officer and Director  
Naval Ship Research & Development Center  
Bethesda, Maryland 20034

Attn: Code 042 (Tech. Lib. Br.)

17 (Struc. Mech. Lab.)

172

172

174

177

1800 (Appl. Math. Lab.)

5412S (Dr. W.D. Sette)

19 (Dr. M.M. Sevik)

1901 (Dr. M. Strassberg)

1945

196 (Dr. D Feit)

1962

Naval Weapons Laboratory  
Dahlgren, Virginia 22448

Naval Research Laboratory  
Washington, D.C. 20375

Attn: Code 8400

8410

8430

8440

6300

6390

6380



Undersea Explosion Research Div.  
Naval Ship R&D Center  
Norfolk Naval Shipyard  
Portsmouth, Virginia 23709  
Attn: Dr. E. Palmer  
Code 780

Naval Ship Research & Development Center  
Annapolis Division  
Annapolis, Maryland 21402  
Attn: Code 2740 - Dr. Y.F. Wang  
28 - Mr. R.J. Wolfe  
281 - Mr. R.B. Niederberger  
2814 - Dr. H. Vanderveldt

Technical Library  
Naval Underwater Weapons Center  
Pasadena Annex  
3202 E. Foothill Blvd.  
Pasadena, California 91107

U.S. Naval Weapons Center  
China Lake, California 93557  
Attn: Code 4062 - Mr. W. Werback  
4520 - Mr. Ken Bischel

Commanding Officer  
U.S. Naval Civil Engr. Lab.  
Code L31  
Port Hueneme, California 93041

Technical Director  
U.S. Naval Ordnance Laboratory  
White Oak  
Silver Spring, Maryland 20910

Technical Director  
Naval Undersea R&D Center  
San Diego, California 92132

Supervisor of Shipbuilding  
U.S. Navy  
Newport News, Virginia 23607

Technical Director  
Mare Island Naval Shipyard  
Vallejo, California 94592

U.S. Navy Underwater Sound Ref. Lab.  
Office of Naval Research  
P.O. Box 8337  
Orlando, Florida 32806

Chief of Naval Operations  
Dept. of the Navy  
Washington, D.C. 20350  
Attn: Code Op077

Strategic Systems Project Office  
Department of the Navy  
Washington, D.C. 20390  
Attn: NSP-001 Chief Scientist

Deep Submergence Systems  
Naval Ship Systems Command  
Code 39522  
Department of the Navy  
Washington, D.C. 20360

Engineering Dept.  
U.S. Naval Academy  
Annapolis, Maryland 21402

Naval Air Systems Command  
Dept. of the Navy  
Washington, D.C. 20360  
Attn: NAVAIR 5302 Aero & Structu  
5308 Structures  
52031F Materials  
604 Tech. Library  
3208 Structures

Director, Aero Mechanics  
Naval Air Development Center  
Johnsville  
Warminster, Pennsylvania 18974

Technical Director  
U.S. Naval Undersea R&D Center  
San Diego, California 92132

Engineering Department  
U.S. Naval Academy  
Annapolis, Maryland 21402

Naval Facilities Engineering Comma  
Dept. of the Navy  
Washington, D.C. 20360  
Attn: NAVFAC 03 Research & Devel  
me

04 " "  
14114 Tech. Library

Naval Sea Systems Command  
Dept. of the Navy  
Washington, D.C. 20360  
Attn: NAVSHIP 03 Res. & Technolog  
031 Ch. Scientist fo  
03412 Hydromechanics  
037 Ship Engineering D

Naval Ship Engineering Center  
Prince George's Plaza  
Hyattsville, Maryland 20782  
Attn: NAVSEC 6100 Ship Sys Engr & Des Dep  
6102C Computer-Aided Ship Des  
6105G  
6110 Ship Concept Design  
6120 Hull Div.  
6120D Hull Div.  
6128 Surface Ship Struct.  
6129 Submarine Struct.

#### Air Force

Commander WADD  
Wright-Patterson Air Force Base  
Dayton, Ohio 45433  
Attn: Code WWRMDD  
AFFDL (FDDS)  
Structures Division  
AFLC (MCEEA)

Chief, Applied Mechanics Group  
U.S. Air Force Inst. of Tech.  
Wright-Patterson Air Force Base  
Dayton, Ohio 45433

Chief, Civil Engineering Branch  
WLRC, Research Division  
Air Force Weapons Laboratory  
Kirtland AFB, New Mexico 87117

Air Force Office of Scientific Research  
1400 Wilson Blvd.  
Arlington, Virginia 22209  
Attn: Mechanics Div.

#### NASA

Structures Research Division  
National Aeronautics & Space Admin.  
Langley Research Center  
Langley Station  
Hampton, Virginia 23365

National Aeronautic & Space Admin.  
Associate Administrator for Advanced  
Research & Technology  
Washington, D.C. 02546

Scientific & Tech. Info. Facility  
NASA Representative (S-AK/DL)  
P.O. Box 5700  
Bethesda, Maryland 20014

#### Other Government Activities

Commandant  
Chief, Testing & Development Div.  
U.S. Coast Guard  
1300 E. Street, N.W.  
Washington, D.C. 20226

Technical Director  
Marine Corps Dev. & Educ. Command  
Quantico, Virginia 22134

Director  
National Bureau of Standards  
Washington, D.C. 20234  
Attn: Mr. B.L. Wilson, EM 219

Dr. M. Gaus  
National Science Foundation  
Engineering Division  
Washington, D.C. 20550

Science & Tech. Division  
Library of Congress  
Washington, D.C. 20540

Director  
Defense Nuclear Agency  
Washington, D.C. 20305  
Attn: SPSS

Commander Field Command  
Defense Nuclear Agency  
Sandia Base  
Albuquerque, New Mexico 87115

Director Defense Research & Engrg  
Technical Library  
Room 3C-128  
The Pentagon  
Washington, D.C. 20301

Chief, Airframe & Equipment Branch  
FS-120  
Office of Flight Standards  
Federal Aviation Agency  
Washington, D.C. 20553

Chief, Research and Development  
Maritime Administration  
Washington, D.C. 20235

Deputy Chief, Office of Ship Constr.  
Maritime Administration  
Washington, D.C. 20235  
Attn: Mr. U.L. Russo

Prof. P.G. Hodge, Jr.  
University of Minnesota  
Dept. of Aerospace Engng & Mechanics  
Minneapolis, Minnesota 55455

Dr. D.C. Drucker  
University of Illinois  
Dean of Engineering  
Urbana, Illinois 61801

Prof. N.M. Newmark  
University of Illinois  
Dept. of Civil Engineering  
Urbana, Illinois 61801

Prof. E. Reissner  
University of California, San Diego  
Dept. of Applied Mechanics  
La Jolla, California 92037

Prof. William A. Nash  
University of Massachusetts  
Dept. of Mechanics & Aerospace Engng.  
Amherst, Massachusetts 01002

Library (Code 0384)  
U.S. Naval Postgraduate School  
Monterey, California 93940

Prof. Arnold Allentuch  
Newark College of Engineering  
Dept. of Mechanical Engineering  
323 High Street  
Newark, New Jersey 07102

Dr. George Herrmann  
Stanford University  
Dept. of Applied Mechanics  
Stanford, California 94305

Prof. J. D. Achenbach  
Northwestern University  
Dept. of Civil Engineering  
Evanston, Illinois 60201

Director, Applied Research Lab.  
Pennsylvania State University  
P. O. Box 30  
State College, Pennsylvania 16801

Prof. Eugen J. Skudrzyk  
Pennsylvania State University  
Applied Research Laboratory  
Dept. of Physics - P.O. Box 30  
State College, Pennsylvania 16801

Prof. J. Kemper  
Polytechnic Institute of Brooklyn  
Dept. of Aero. Engrg. & Applied Mech.  
333 Jay Street  
Brooklyn, N.Y. 11201

Prof. J. Klosner  
Polytechnic Institute of Brooklyn  
Dept. of Aerospace & Appl. Mech.  
333 Jay Street  
Brooklyn, N.Y. 11201

Prof. R.A. Schapery  
Texas A&M University  
Dept. of Civil Engineering  
College Station, Texas 77840

Prof. W.D. Pilkey  
University of Virginia  
Dept. of Aerospace Engineering  
Charlottesville, Virginia 22903

Dr. H.G. Schaeffer  
University of Maryland  
Aerospace Engineering Dept.  
College Park, Maryland 20742

Prof. K.D. Willmert  
Clarkson College of Technology  
Dept. of Mechanical Engineering  
Potsdam, N.Y. 13676

Dr. J.A. Stricklin  
Texas A&M University  
Aerospace Engineering Dept.  
College Station, Texas 77843

Dr. L.A. Schmit  
University of California, LA  
School of Engineering & Applied Science  
Los Angeles, California 90024

Dr. H.A. Kamel  
The University of Arizona  
Aerospace & Mech. Engineering Dept.  
Tucson, Arizona 85721

Dr. B.S. Berger  
University of Maryland  
Dept. of Mechanical Engineering  
College Park, Maryland 20742

Prof. G. R. Irwin  
Dept. of Mechanical Engrg.  
University of Maryland  
College Park, Maryland 20742

Atomic Energy Commission  
Div. of Research & Development  
Germantown, Maryland 20767

Ship Hull Research Committee  
National Research Council  
National Academy of Sciences  
2101 Constitution Avenue  
Washington, D.C. 20418  
Attn: Mr. A.R. Lytle

PART 2 - CONTRACTORS AND OTHER  
TECHNICAL COLLABORATORS

Universities

Dr. J. Tinsley Oden  
University of Texas at Austin  
345 Eng. Science Bldg.  
Austin, Texas 78712

Prof. Julius Miklowitz  
California Institute of Technology  
Div. of Engineering & Applied Sciences  
Pasadena, California 91109

Dr. Harold Liebowitz, Dean  
School of Engr. & Applied Science  
George Washington University  
725 - 23rd St., N.W.  
Washington, D.C. 20006

Prof. Eli Sternberg  
California Institute of Technology  
Div. of Engr. & Applied Sciences  
Pasadena, California 91109

Prof. Paul M. Naghdi  
University of California  
Div. of Applied Mechanics  
Etcheverry Hall  
Berkeley, California 94720

Professor P. S. Symonds  
Brown University  
Division of Engineering  
Providence, R.I. 02912

Prof. A. J. Durelli  
The Catholic University of America  
Civil/Mechanical Engineering  
Washington, D.C. 20017

Prof. R.D. Testa  
Columbia University  
Dept. of Civil Engineering  
S.W. Mudd Bldg.  
New York, N.Y. 10027

Prof. H. H. Bleich  
Columbia University  
Dept. of Civil Engineering  
Amsterdam & 120th St.  
New York, N.Y. 10027

Prof. F.L. DiMaggio  
Columbia University  
Dept. of Civil Engineering  
616 Mudd Building  
New York, N.Y. 10027

Prof. A.M. Freudenthal  
George Washington University  
School of Engineering &  
Applied Science  
Washington, D.C. 20006

D. C. Evans  
University of Utah  
Computer Science Division  
Salt Lake City, Wash 84112

Prof. Norman Jones  
Massachusetts Inst. of Technology  
Dept. of Naval Architecture &  
Marine Engrng  
Cambridge, Massachusetts 02139

Professor Albert I. King  
Biomechanics Research Center  
Wayne State University  
Detroit, Michigan 48202

Dr. V. R. Hodgson  
Wayne State University  
School of Medicine  
Detroit, Michigan 48202

Dean B. A. Boley  
Northwestern University  
Technological Institute  
2145 Sheridan Road  
Evanston, Illinois 60201

### Industry and Research Institutes

Library Services Department  
Report Section Bldg. 14-14  
Argonne National Laboratory  
9700 S. Cass Avenue  
Argonne, Illinois 60440

Dr. M. C. Junger  
Cambridge Acoustical Associates  
129 Mount Auburn St.  
Cambridge, Massachusetts 02138

Dr. L.H. Chen  
General Dynamics Corporation  
Electric Boat Division  
Groton, Connecticut 06340

Dr. J.E. Greenspon  
J.G. Engineering Research Associates  
3831 Menio Drive  
Baltimore, Maryland 21215

Dr. S. Batdorf  
The Aerospace Corp.  
P.O. Box 92957  
Los Angeles, California 90009

Dr. K.C. Park  
Lockheed Palo Alto Research Laboratory  
Dept. 5233, Bldg. 205  
3251 Hanover Street  
Palo Alto, CA 94304

Library  
Newport News Shipbuilding &  
Dry Dock Company  
Newport News, Virginia 23607

Dr. W.F. Bozich  
McDonnell Douglas Corporation  
5301 Bolsa Ave.  
Huntington Beach, CA 92647

Dr. H.N. Abramson  
Southwest Research Institute  
Technical Vice President  
Mechanical Sciences  
P.O. Drawer 28510  
San Antonio, Texas 78284

Dr. R.C. DeHart  
Southwest Research Institute  
Dept. of Structural Research  
P.O. Drawer 28510  
San Antonio, Texas 78284

Dr. M.L. Baron  
Weidlinger Associates,  
Consulting Engineers  
110 East 59th Street  
New York, N.Y. 10022

Dr. W.A. von Rieseemann  
Sandia Laboratories  
Sandia Base  
Albuquerque, New Mexico 87115

Dr. T.L. Geers  
Lockheed Missiles & Space Co.  
Palo Alto Research Laboratory  
3251 Hanover Street  
Palo Alto, California 94304

Dr. J.L. Tocher  
Boeing Computer Services, Inc.  
P.O. Box 24346  
Seattle, Washington 98124

Mr. William Caywood  
Code BBE, Applied Physics Laboratory  
8621 Georgia Avenue  
Silver Spring, Maryland 20034

Mr. P.C. Durup  
Lockheed-California Company  
Aeromechanics Dept., 74-43  
Burbank, California 91503

Dr. S.J. Fenves  
Carnegie-Mellon University  
Dept. of Civil Engineering  
Schenley Park  
Pittsburgh, Pennsylvania 15213

Dr. Ronald L. Huston  
Dept. of Engineering Analysis  
Mail Box 112  
University of Cincinnati  
Cincinnati, Ohio 45221

Prof. George Sih  
Dept. of Mechanics  
Lehigh University  
Bethlehem, Pennsylvania 18015

Prof. A.S. Kobayashi  
University of Washington  
Dept. of Mechanical Engineering  
Seattle, Washington 98105

Librarian  
Webb Institute of Naval Architecture  
Crescent Beach Road, Glen Cove  
Long Island, New York 11542

Prof. Daniel Frederick  
Virginia Polytechnic Institute  
Dept. of Engineering Mechanics  
Blacksburg, Virginia 24061

Prof. A.C. Eringen  
Dept. of Aerospace & Mech. Sciences  
Princeton University  
Princeton, New Jersey 08540

Dr. S.L. Koh  
School of Aero., Astro. & Engr. Sc.  
Purdue University  
Lafayette, Indiana 47907

Prof. E.H. Lee  
Div. of Engrg. Mechanics  
Stanford University  
Stanford, California 94305

Prof. R.D. Mindlin  
Dept. of Civil Engrg.  
Columbia University  
S.W. Mudd Building  
New York, N.Y. 10027

Prof. S.B. Dong  
University of California  
Dept. of Mechanics  
Los Angeles, California 90024  
Prof. Burt Paul  
University of Pennsylvania  
Towne School of Civil & Mech. Engrg.  
Rm. 113 - Towne Building  
220 S. 33rd Street  
Philadelphia, Pennsylvania 19104  
Prof. H.W. Liu  
Dept. of Chemical Engr. & Metal.  
Syracuse University  
Syracuse, N.Y. 13210

Prof. S. Bodner  
Technion R&D Foundation  
Haifa, Israel

Prof. R.J.H. Bollard  
Chairman, Aeronautical Engr. Dept.  
207 Guggenheim Hall  
University of Washington  
Seattle, Washington 98105

Prof. G.S. Heller  
Division of Engineering  
Brown University  
Providence, Rhode Island 02912

Prof. Werner Goldsmith  
Dept. of Mechanical Engineering  
Div. of Applied Mechanics  
University of California  
Berkeley, California 94720

Prof. J.R. Rice  
Division of Engineering  
Brown University  
Providence, Rhode Island 02912

Prof. R.S. Rivlin  
Center for the Application of Mathem.  
Lehigh University  
Bethlehem, Pennsylvania 18015

Library (Code 0384)  
U.S. Naval Postgraduate School  
Monterey, California 93940

Dr. Francis Cozzarelli  
Div. of Interdisciplinary  
Studies & Research  
School of Engineering  
State University of New York  
Buffalo, N.Y. 14214